

## ON A SOLVABILITY OF THE NONLINEAR INVERSE BOUNDARY VALUE PROBLEM FOR THE BOUSSINESQ EQUATION

A.S. Farajov\*

Azerbaijan State Pedagogical University, Baku, Azerbaijan

---

**Abstract.** We study the classical solution of the nonlinear inverse boundary value problem for the sixth-order Boussinesq equation with double dispersion. The essence of the problem is that it is required together with the solution to determine the unknown coefficient. The problem is considered in a rectangular area. To solve the considered problem, the transition from the original inverse problem to some auxiliary inverse problem is carried out. The existence and uniqueness of a solution to the auxiliary problem are proved with the help of contracted mappings. Then the transition to the original inverse problem is made, as a result, a conclusion is made about the solvability of the original inverse problem.

---

**Keywords:** inverse boundary value problem, classical solution, uniqueness, existence, Fourier method, Boussinesq equation.

**AMS Subject Classification:** 35R30.

**Corresponding author:** A.S. Farajov, Azerbaijan State Pedagogical University, U. Hajibeyli Str. 68, AZ1000, Baku, Azerbaijan, e-mail: [a.farajov@mail.ru](mailto:a.farajov@mail.ru)

*Received: 28 February 2022; Revised: 16 April 2022; Accepted: 3 June 2022; Published: 5 August 2022.*

---

## 1 Introduction

The theory of inverse problems for the differential equations is a dynamically developing branch of mathematics. Inverse problems arise in various areas of human activity such as seismology, mineral exploration, biology, medicine, quality control of industrial products, etc., which puts them among the topical problems. The presence of additional unknown functions in the inverse problems requires some additional redefinition conditions are also given. The fundamental works in studying the inverse problems both from theoretical and practical points of view belong to outstanding scientists A.N. Tikhonov, M.M. Lavrentiev, V.K. Ivanov and their students.

The sixth-order Boussinesq equation with double dispersion describes motion of waves on water with a stress surface and was considered by Schneider and Eugene in (Ivanov, 1962). Various boundary value problems for the Boussinesq type equation were studied in (Polat, 2008; Xu et al., 2008; Polat & Ertas, 2009; Lin et al., 2009; Wang, 2016; Wang & Esfahani, 2014). Various inverse problems for certain types of partial differential equations have been studied in many works (Kamynin, 2013; 2018; 2020). For the Boussinesq equation of the sixth order, inverse problems were considered in (Farajov, 2021a; 2021b; 2021c; Mehraliyev & Farajov, 2021).

In this paper we consider the inverse problem for the sixth-order Boussinesq equation with double dispersion, where, along with the solution, an unknown coefficient should also be found.

## 2 Formulation of the problem and its equivalent form

Let  $D_T = \{(x, t) : 0 \leq x \leq 1, 0 \leq t \leq T\}$  and  $f(x, t)$ ,  $\varphi(x)$ ,  $\psi(x)$ ,  $h(t)$  are given functions defined for  $x \in [0, 1]$ ,  $t \in [0, T]$ . Consider the following inverse problem: to find a pair

$\{u(x, t), a(t)\}$  of the functions  $u(x, t), a(t)$  satisfying the equation

$$\begin{aligned} u_{tt}(x, t) - u_{xx}(x, t) - u_{ttxx}(x, t) + u_{xxxx}(x, t) + u_{ttxxxx}(x, t) = \\ = a(t)u(x, t) + f(x, t) \end{aligned} \tag{1}$$

with initial

$$u(x, 0) = \varphi(x), \quad u_t(x, 0) = \psi(x) \quad (0 \leq x \leq 1) \tag{2}$$

and boundary conditions

$$u_x(0, t) = u(1, t) = u_{xxx}(0, t) = u_{xx}(1, t) = 0 \quad (0 \leq t \leq T) \tag{3}$$

and with additional condition

$$u(0, t) = h(t) \quad (0 \leq t \leq T), \tag{4}$$

Introduce the designation

$$\tilde{C}^{4,2}(D_T) = \{u(x, t) : u(x, t) \in C^2(D_T), u_{xxxx}(x, t), u_{ttxxxx}(x, t) \in C(D_T)\}$$

**Definition 1.** A pair  $\{u(x, t), a(t)\}$  of the functions  $u(x, t) \in C^{4,2}(D_T)$  and  $a(t) \in C[0, T]$  satisfying equation (1) in  $D_T$ , condition (2) in  $[0, 1]$  and conditions (3)-(4) in  $[0, T]$  we call a classical solution to boundary value (1)-(4).

We prove the following

**Theorem 1.** Let  $f(x, t) \in C(D_T)$ ,  $\varphi(x), \psi(x) \in C[0, 1]$ ,  $h(t) \in C^2[0, T]$ ,  $h(t) \neq 0$  ( $0 \leq t \leq T$ ) and the matching conditions

$$\varphi(0) = h(0), \quad \psi(0) = h'(0) .$$

are satisfied. Then the problem of finding a classical solution to problem (1)-(4) is equivalent to the problem of determining the functions  $u(x, t) \in C^{4,2}(D_T)$  and  $a(t) \in C[0, T]$  from (1)-(3) and

$$\begin{aligned} h''(t) - u_{xx}(0, t) - u_{ttxx}(0, t) + u_{xxxx}(0, t) + \\ + u_{ttxxxx}(0, t) = a(t)h(t) + f(0, t) \quad (0 \leq t \leq T) . \end{aligned} \tag{5}$$

*Proof.* Let  $\{u(x, t), a(t)\}$  be a classical solution to problem (1)-(4). Since  $h(t) \in C^2[0, T]$ , differentiating (4) two times over  $t$  we get

$$u_t(0, t) = h'(t), \quad u_{tt}(0, t) = h''(t) \quad (0 \leq t \leq T). \tag{6}$$

Taking  $x = 0$  in equation (1) we find

$$\begin{aligned} u_{tt}(0, t) - u_{xx}(0, t) - u_{ttxx}(0, t) + u_{xxxx}(0, t) + u_{ttxxxx}(0, t) = \\ = a(t)u(0, t) + f(0, t) \quad (0 \leq t \leq T). \end{aligned} \tag{7}$$

From this considering (4) and (6) we arrive at (5).

Now let's suppose that  $\{u(x, t), a(t)\}$  is a solution of problem (1)-(3), (5). Then from (5) and (7) we get

$$\frac{d^2}{dt^2}(u(0, t) - h(t)) = a(t)(u(0, t) - h(t)) \quad (0 \leq t \leq T) \tag{8}$$

Considering (2) and  $\varphi(0) = h(0), \psi(0) = h'(0)$  we have

$$u_t(x, 0) - h'(0) = \psi(0) - h'(0) = 0 . \tag{9}$$

From (8), taking into account (9), it is clear that condition (4) is also satisfied. The theorem is proved.  $\square$

### 3 Solvability of the inverse boundary value problem

The first component  $u(x, t)$  of the solution  $\{u(x, t), a(t)\}$  to problem (1)-(3), (5) we seek in the form

$$u(x, t) = \sum_{k=1}^{\infty} u_k(t) \cos \lambda_k x \quad \left( \lambda_k = \frac{\pi}{2}(2k-1) \right), \quad (10)$$

where

$$u_k(t) = 2 \int_0^1 u(x, t) \cos \lambda_k x dx \quad (k = 1, 2, \dots).$$

Then applying the formal Fourier scheme, from (1) and (2) we obtain

$$(1 + \lambda_k^2 + \lambda_k^4)u_k''(t) + (\lambda_k^2 + \lambda_k^4)u_k(t) = F_k(t; u, a) \quad (0 \leq t \leq T; k = 1, 2, \dots), \quad (11)$$

$$u_k(0) = \varphi_k, \quad u_k'(0) = \psi_k \quad (k = 1, 2, \dots), \quad (12)$$

where

$$F_k(t; u, a) = a(t)u_k(t) + f_k(t), \quad f_k(t) = 2 \int_0^1 f(x, t) \cos \lambda_k x dx,$$

$$\varphi_k = 2 \int_0^1 \varphi(x) \cos \lambda_k x dx, \quad \psi_k = 2 \int_0^1 \psi(x) \cos \lambda_k x dx \quad (k = 1, 2, \dots).$$

Solving problem (11)-(12) we find

$$u_k(t) = \varphi_k \cos \beta_k t + \frac{1}{\beta_k} \psi_k \sin \beta_k t +$$

$$+ \frac{1}{\beta_k(1 + \lambda_k^2 + \lambda_k^4)} \int_0^t F_k(\tau; u, a) \sin \beta_k(t - \tau) d\tau \quad (k = 1, 2, \dots), \quad (13)$$

where

$$\beta_k^2 = \frac{\lambda_k^2 + \lambda_k^4}{1 + \lambda_k^2 + \lambda_k^4} \quad (k = 1, 2, \dots).$$

After substitution of the expression  $u_k(t)$  ( $k = 1, 2, \dots$ ) into (10) for the determination of  $u(x, t)$  we get

$$u(x, t) = \sum_{k=1}^{\infty} \left\{ \varphi_k \cos \beta_k t + \frac{1}{\beta_k} \psi_k \sin \beta_k t + \right.$$

$$\left. + \frac{1}{\beta_k(1 + \lambda_k^2 + \lambda_k^4)} \int_0^t F_k(\tau; u, a) \sin \beta_k(t - \tau) d\tau \right\} \cos \lambda_k x. \quad (14)$$

Now from (5) taking into account (10) we have

$$a(t) = [h(t)]^{-1} \left\{ h''(t) - f(0, t) + \sum_{k=1}^{\infty} (\lambda_k^2 + \lambda_k^4)(u_k''(t) + u_k(t))u_k(t) \right\}. \quad (15)$$

Consideration of (13) in (11) gives

$$(\lambda_k^2 + \lambda_k^4)(u_k''(t) + u_k(t)) = -u_k''(t) + F_k(t; u, a, b) =$$

$$= \frac{\lambda_k^2 + \lambda_k^4}{1 + \lambda_k^2 + \lambda_k^4} u_k(t) + \left( 1 - \frac{1}{1 + \lambda_k^2 + \lambda_k^4} \right) F_k(t; u, a, b) =$$

$$= \frac{\lambda_k^2 + \lambda_k^4}{1 + \lambda_k^2 + \lambda_k^4} u_k(t) + \frac{\lambda_k^2 + \lambda_k^4}{1 + \lambda_k^2 + \lambda_k^4} F_k(t; u, a, b) = \beta_k^2 u_k(t) + \beta_k^2 F_k(t; u, a, b) =$$

$$\begin{aligned}
 &= \beta_k^2 u_k(t) + \beta_k^2 F_k(t; u, a, b) = \beta_k^2 \left[ \varphi_k \cos \beta_k t + \frac{1}{\beta_k} \psi_k \sin \beta_k t + \right. \\
 &+ \left. \frac{1}{\beta_k(1 + \lambda_k^2 + \lambda_k^4)} \int_0^t F_k(\tau; u, a, b) \sin \beta_k(t - \tau) d\tau \right] + \beta_k^2 F_k(t; u, a, b), \\
 &k = 1, 2, \dots, \quad 0 \leq t \leq T.
 \end{aligned}$$

To obtain an equation for the second component  $a(t)$  of the solution  $\{u(x, t), a(t)\}$  we put the last relation into (15)

$$\begin{aligned}
 a(t) = [h(t)]^{-1} &\left\{ h''(t) - f(0, t) + \sum_{k=1}^{\infty} \beta_k^2 \left[ \varphi_k \cos \beta_k t + \frac{1}{\beta_k} \psi_k \sin \beta_k t + \right. \right. \\
 &+ \left. \left. \frac{1}{\beta_k(1 + \lambda_k^2 + \lambda_k^4)} \int_0^t F_k(\tau; u, p) \sin \beta_k(t - \tau) d\tau + F_k(t; u, a, b) \right] \right\}. \tag{16}
 \end{aligned}$$

Thus, solution of problem (1)-(3), (5) is reduced to the solution of system (14), (16) with respect to the unknown functions  $u(x, t)$  and  $a(t)$ .

To study the problem of the uniqueness of the solution of problem (1)-(3), (5), the following lemma plays an important role.

**Lemma 1.** *If  $\{u(x, t), a(t)\}$  is arbitrary classical solution of problem (1)-(3), (5), then the function*

$$u_k(t) = 2 \int_0^1 u(x, t) \cos \lambda_k x dx \quad (k = 1, 2, \dots)$$

*satisfies system (13) in  $[0, T]$ .*

*Proof.* Let  $\{u(x, t), a(t)\}$  be any solution to problem (1)-(3), (5). Then multiplying both sides of equation (1) by the function  $2 \cos \lambda_k x$  ( $k = 1, 2, \dots$ ), integrating the obtained equality over  $x$  from 0 to 1 and using the relations

$$\begin{aligned}
 2 \int_0^1 u_{tt}(x, t) \cos \lambda_k x dx &= \frac{d^2}{dt^2} \left( 2 \int_0^1 u(x, t) \cos \lambda_k x dx \right) = u_k''(t) \quad (k = 1, 2, \dots), \\
 2 \int_0^1 u_{xx}(x, t) \cos \lambda_k x dx &= -\lambda_k^2 \left( 2 \int_0^1 u(x, t) \cos \lambda_k x dx \right) = -\lambda_k^2 u_k(t) \quad (k = 1, 2, \dots), \\
 2 \int_0^1 u_{ttxx}(x, t) \cos \lambda_k x dx &= -\lambda_k^2 \left( 2 \int_0^1 u_{tt}(x, t) \cos \lambda_k x dx \right) = -\lambda_k^2 u_k''(t) \quad (k = 1, 2, \dots) \\
 2 \int_0^1 u_{xxxx}(x, t) \cos \lambda_k x dx &= \lambda_k^4 \left( 2 \int_0^1 u(x, t) \cos \lambda_k x dx \right) = \lambda_k^4 u_k(t) \quad (k = 1, 2, \dots), \\
 2 \int_0^1 u_{ttxxxx}(x, t) \cos \lambda_k x dx &= \lambda_k^4 \left( 2 \int_0^1 u_{tt}(x, t) \cos \lambda_k x dx \right) = -\lambda_k^4 u_k(t) \quad (k = 1, 2, \dots)
 \end{aligned}$$

we obtain that equation (11) is satisfied.

Similarly, the fulfilment of (12) is obtained from (2). Thus  $u_k(t)$  ( $k = 1, 2, \dots$ ) is a solution to problem (11), (12).

As immediately follows from this the function  $u_k(t)$  ( $k = 1, 2, \dots$ ) satisfies to system (13) on  $[0, T]$ . Lemma is proved.  $\square$

It is obvious that if  $u_k(t) = 2 \int_0^1 u(x, t) \cos \lambda_k x dx$  ( $k = 1, 2, \dots$ ) is a solution of system (13), then the pair  $\{u(x, t), a(t)\}$  of the functions  $u(x, t) = \sum_{k=1}^{\infty} u_k(t) \cos \lambda_k x$  and  $a(t)$  is a solution to system (14), (16).

The above lemma implies the validity of the following

**Consequence.** Let system (14), (16) have a unique solution. Then problem (1)-(3), (5) cannot have more than one solution, i.e. if problem (1)-(3), (5) has a solution, then it is unique.

Now, in order to study problem (1)-(3), (5) consider the following spaces.

1. Denote by  $B_{2,T}^5$  [19] the set of all functions  $u(x, t)$  of the form

$$u(x, t) = \sum_{k=1}^{\infty} u_k(t) \cos \lambda_k x \quad \left( \lambda_k = \frac{\pi}{2}(2k - 1) \right),$$

Defined on  $D_T$ , where each of the functions  $u_k(t)$  ( $k = 1, 2, \dots$ ) is continuous on  $[0, T]$  and

$$J_T(u) \equiv \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|u_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} < +\infty.$$

The norm in this space is defined as

$$\|u(x, t)\|_{B_{2,T}^5} = J(u).$$

2. By  $E_T^5$  we denote the space of the vector functions  $\{u(x, t), a(t)\}$  such that  $u(x, t) \in B_{2,T}^5$ ,  $a(t) \in C[0, T]$  and equip this space by the norm

$$\|z\|_{E_T^5} = \|u(x, t)\|_{B_{2,T}^5} + \|a(t)\|_{C[0,T]}.$$

Clearly,  $B_{2,T}^5$  and  $E_T^5$  are Banach spaces.

Now we consider in  $E_T^5$  the operator

$$\Phi(u, a) = \{\Phi_1(u, a), \Phi_2(u, a)\},$$

where

$$\Phi_1(u, a) = \tilde{u}(x, t) \equiv \sum_{k=1}^{\infty} \tilde{u}_k(t) \cos \lambda_k x, \quad \Phi_2(u, a) = \tilde{a}(t),$$

$\tilde{u}_k(t)$  ( $k = 1, 2, \dots$ ) and  $\tilde{a}(t)$  are the right hand sides of (13) and (16), correspondingly.

Obviously

$$\frac{\sqrt{3}}{3} < \beta_k < \sqrt{2}, \quad \frac{\sqrt{2}}{2} < \frac{1}{\beta_k} < \sqrt{3}.$$

Then we have

$$\begin{aligned} & \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|\tilde{u}_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} \leq 2 \left( \sum_{k=1}^{\infty} (\lambda_k^5 |\varphi_k|)^2 \right)^{\frac{1}{2}} + 2\sqrt{3} \left( 5 \sum_{k=1}^{\infty} (\lambda_k^5 |\psi_k|)^2 \right)^{\frac{1}{2}} + \\ & + 2\sqrt{3T} \left( \int_0^T \sum_{k=1}^{\infty} (\lambda_k |f_k(\tau)|)^2 d\tau \right)^{\frac{1}{2}} + 2\sqrt{3} T \|a(t)\|_{C[0,T]} \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|u_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}}, \quad (17) \\ & \|\tilde{a}(t)\|_{C[0,T]} = \left\| [h(t)]^{-1} \right\|_{C[0,T]} \left\{ \|h''(t) - f(0, t)\|_{C[0,T]} + \right. \\ & \left. + \sqrt{2} \left( \sum_{k=1}^{\infty} \lambda_k^{-2} \right)^{\frac{1}{2}} \left[ \left( \sum_{k=1}^{\infty} (\lambda_k^5 |\varphi_k|)^2 \right)^{\frac{1}{2}} + \sqrt{3} \left( \sum_{k=1}^{\infty} (\lambda_k^5 |\psi_k|)^2 \right)^{\frac{1}{2}} \right] \right\} \end{aligned}$$

$$\begin{aligned}
 & +\sqrt{3T} \left( \int_0^T \sum_{k=1}^{\infty} (\lambda_k^4 |f_k(\tau)|)^2 d\tau \right)^{\frac{1}{2}} + \sqrt{3} T \|a(t)\|_{C[0,T]} \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|u_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} \\
 & + \left( \sum_{k=1}^{\infty} (\lambda_k \|f_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \|a(t)\|_{C[0,T]} \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|u_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} \Bigg\} \quad (18)
 \end{aligned}$$

Assume that the data of problem (1)-(3), (5) satisfy the following conditions: ]]

1.  $\varphi(x) \in C^4[0, 1]$ ,  $\varphi^{(5)}(x) \in L_2(0, 1)$ ,  $\varphi'(0) = \varphi(1) = \varphi'''(0) = \varphi''(1) = \varphi^{(4)}(1) = 0'$ .
2.  $\psi(x) \in C^4[0, 1]$ ,  $\psi^{(5)}(x) \in L_2(0, 1)$ ,  $\psi'(0) = \psi(1) = \psi'''(0) = \psi''(1) = \psi^{(4)}(1) = 0$ .
3.  $f(x, t) \in C(D_T)$ ,  $f_x(x, t) \in L_2(D_T)$ ,  $f(1, t) = 0$  ( $0 \leq t \leq T$ ).
4.  $h(t) \in C^2[0, T]$ ,  $h(t) \neq 0$  ( $0 \leq t \leq T$ ).

Then from (17)-(18) we have

$$\|\tilde{u}(x, t)\|_{B_{2,T}^7} \leq A_1(T) + B_1(T) \|a(t)\|_{C[0,T]} \|u(x, t)\|_{B_{2,T}^7}, \quad (19)$$

$$\|\tilde{a}(t)\|_{C[0,T]} \leq A_2(T) + B_2(T) \|a(t)\|_{C[0,T]} \|u(x, t)\|_{B_{2,T}^7}, \quad (20)$$

where

$$A_1(T) = 2 \left\| \varphi^{(5)}(x) \right\|_{L_2(0,1)} + 2\sqrt{3} \left\| \psi^{(5)}(x) \right\|_{L_2(0,1)} + 2\sqrt{3T} 2 \|f_x(x, t)\|_{L_2(D_T)},$$

$$B_1(T) = 2\sqrt{3}T,$$

$$\begin{aligned}
 A_2(T) = & \left\| [h(t)]^{-1} \right\|_{C[0,T]} \left\{ \|h''(t) - f(0, t)\|_{C[0,T]} + \left( \sum_{k=1}^{\infty} \lambda_k^{-2} \right)^{\frac{1}{2}} \times \times \left[ \left\| \varphi^{(5)}(x) \right\|_{L_2(0,1)} + \right. \right. \\
 & \left. \left. + \sqrt{3} \left\| \psi^{(3)}(x) \right\|_{L_2(0,1)} + \sqrt{3T} \|f_x(x, t)\|_{L_2(D_T)} \left\| \|f_x(x, t)\|_{C[0,T]} \right\|_{L_2(0,1)} \right] \right\},
 \end{aligned}$$

$$B_2(T) = \left\| [h(t)]^{-1} \right\|_{C[0,T]} \left( \sum_{k=1}^{\infty} \lambda_k^{-2} \right)^{\frac{1}{2}} (T + 1).$$

From inequalities (19)-(20) we conclude

$$\|\tilde{u}(x, t)\|_{B_{2,T}^5} + \|\tilde{a}(t)\|_{C[0,T]} \leq A(T) + B(T) \|a(t)\|_{C[0,T]} \|u(x, t)\|_{B_{2,T}^5}, \quad (21)$$

where

$$A(T) = A_1(T) + A_2(T), B(T) = B_1(T) + B_2(T).$$

So, we can prove the following theorem:

**Theorem 2.** *Let conditions 1-4 be satisfied and*

$$(A(T) + 2)^2 B(T) < 1. \quad (22)$$

*The problem (1)-(3), (5) has a unique solution in the ball  $K = K_R(\|z\|_{E_T^5} \leq R = A(T) + 2)$  of the space  $E_T^5$ .*

*Proof.* In the space  $E_T^5$  consider the equation

$$z = \Phi z, \tag{23}$$

where  $z = \{u, a\}$ , the components  $\Phi_i(u, a) (i = 1, 2)$  of the operator  $\Phi(u, a)$  are defined by the right hand sides of equations (14) and (16) .

Consider the operator  $\Phi(u, a)$  in the ball  $K = K_R$  from  $E_T^5$ . Similarly to (22) we obtain that the estimations

$$\|\Phi z\|_{E_T^5} \leq A(T) + B(T) \|a(t)\|_{C[0,T]} \|u(x, t)\|_{B_{2,T}^7}, \tag{24}$$

$$\|\Phi z_1 - \Phi z_2\|_{E_T^5} \leq$$

$$\leq B(T)R \left( \|a_1(t) - a_2(t)\|_{C[0,T]} + \|u_1(x, t) - u_2(x, t)\|_{B_{2,T}^5} \right) \tag{25}$$

for the arbitrary  $z, z_1, z_2 \in K_R$  . Then, from estimates (24), (25), taking into account (22), it follows that the operator  $\Phi$  acts in the ball and is contractive. Therefore in the ball  $K = K_R$  the operator  $\Phi$  has a single fixed point  $\{u, a\}$  which is a unique solution to equation (23) in the ball  $K = K_R$ , i.e.  $\{u, a\}$  is a unique solution to system (14)-(16) in the ball  $K = K_R$ .

The function  $u(x, t)$  as an element of the space  $B_{2,T}^5$  has continuous derivatives  $u(x, t), u_x(x, t), u_{xx}(x, t), u_{xxx}(x, t), u_{xxxx}(x, t)$  in  $D_T$ .

As one can easily see from

$$\begin{aligned} \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|u_k''(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} &\leq \sqrt{2} \left( \sum_{k=1}^{\infty} (\lambda_k^5 \|u_k(t)\|_{C[0,T]})^2 \right)^{\frac{1}{2}} + \\ &+ \sqrt{2} \left\| \|f_x(x, t) + a(t)u_x(x, t) + b(t)u_{tx}(x, t)\|_{C[0,T]} \right\|_{L_2(0,1)}. \end{aligned}$$

It implies that  $u_{tt}(x, t), u_{ttx}(x, t), u_{ttxx}(x, t), u_{ttxxx}(x, t), u_{ttxxxx}(x, t)$  are continuous in  $D_T$ .

It is easy to check that equation (1) and conditions (2), (3) and (5) are satisfied in the usual sense. Therefore,  $\{u(x, t), a(t)\}$  is a solution to problem (1)-(3), (5), and, by virtue of the corollary of Lemma 1, it is unique in the ball  $K = K_R$ . The theorem is proved.  $\square$

Using Theorem 1, we prove the following

**Theorem 3.** *Let all conditions of Theorem 2 be satisfied and*

$$\varphi(0) = h(0), \quad \psi(0) = h'(0) .$$

*The problem (1)-(4) has unique classical solution in the ball  $K = K_R \left( \|z\|_{E_T^5} \leq R = A(T) + 2 \right)$ .*

## References

- Abdullaev, V.M., Aida-Zade K.R. (2016). Finite-difference methods for solving loaded parabolic equations. *Computational Mathematics and Mathematical Physics*, 56(1), 93-105.
- Farajov, A.S. (2021a). Inverse boundary value problem for the sixth-order Boussinesq equation with double variance, *News of Baku University, Series of physico-mathematical sciences*, 3, 16-27.
- Farajov, A.S. (2021b). On a non-local boundary value problem for the sixth-order Boussinesq equation with double variance. *Transactions of Azerbaijan Ped. Univ., Ser. of math. and natural sciences*, 69(2), 22-33.

- Farajov, A.S. (2021c). On a non-local inverse boundary value problem for the Boussinesq equation. *International Conference on Differential Equations, Mathematical Modeling and Computational Algorithms*: Belgorod, October 25-29, 243-24.
- Ivanov V.K. (1962). On linear incorrect problems. *Doklady USSR Academy of Sciences*, 145(2), 270-272.
- Kamynin, V. L. (2020). The inverse problem of simultaneous determination of the two time-dependent lower coefficients in a nondivergent parabolic equation in the plane. *Mathematical Notes*, 107(1), 93-104.
- Kamynin, V.L. (2013). The inverse problem of determining the lower-order coefficient in parabolic equations with integral observation. *Mathematical Notes*, 94(1), 205-213.
- Kamynin, V.L. (2018). On inverse problems for strongly degenerate parabolic equations under the integral observation condition. *Computational Mathematics and Mathematical Physics*, 58(12), 2002-2017.
- Khudaverdiev K.I., Veliyev A.A. (2010). Study of a one-dimensional mixed problem for a class of third-order pseudohyperbolic equations with a nonlinear operator right-hand side. Baku, 2010, 168 p.
- Lavrent'ev M.M. (1962). On some incorrect problems of mathematical physics. Novosibirsk, Nauka.
- Lin, Q., Wu, Y. H., & Loxton, R. (2009). On the Cauchy problem for a generalized Boussinesq equation. *Journal of Mathematical Analysis and Applications*, 353(1), 186-195.
- Mehraliyev, Y.T., Farajov, A.S. (2021). On solvability of an inverse boundary value problem for double dispersive sixth order Boussinesq equation. Fundamental and applied problems mathematics and computer science. Materials of the XIV International Conference, Makhachkala, September, 16-19, 157-160.
- Polat, N. (2008). Existence and blow up of solutions of the Cauchy problem of the generalized damped multidimensional improved modified Boussinesq equation. *Zeitschrift Für Naturforschung A*, 63(9), 543-552.
- Polat, N., Ertaş, A. (2009). Existence and blow-up of solution of Cauchy problem for the generalized damped multidimensional Boussinesq equation. *Journal of Mathematical Analysis and Applications*, 349(1), 10-20.
- Schneider, G., Eugene, C.W. (2001). Kawahara dynamics in dispersive media. *Physica D: Nonlinear Phenomena*, 152, 384-394.
- Tikhonov A.N. (1943). On the stability of inverse problems. *Doklady USSR Academy of Sciences*, 39(5), 195-198.
- Wang, H., Esfahani, A. (2014). Global rough solutions to the sixth-order Boussinesq equation. *Nonlinear Analysis: Theory, Methods & Applications*, 102, 97-104.
- Wang, Y. (2016). Cauchy problem for the sixth-order damped multidimensional Boussinesq equation. *Electronic Journal of Differential Equations*, 64, 1-16.
- Xu, R., Liu, Y., & Liu, B. (2011). The Cauchy problem for a class of the multidimensional Boussinesq-type equation. *Nonlinear Analysis: Theory, Methods & Applications*, 74(6), 2425-2437.